

**Solution Manual for Electronics A Physical Approach 1st Edition by Snoke
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Chapter 2 Exercises

2.1

The electric field outside a charged sphere is the same as for a point source,

$$E(r) = \frac{Q}{4\pi\epsilon_0 r^2};$$

where Q is the charge on the inner surface of radius a. The potential drop is the integral

$$V = \int_a^b \frac{Q}{4\pi\epsilon_0 r^2} dr = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{a} - \frac{1}{b} \right];$$

The capacitance is therefore

$$C = \frac{Q}{V} = 4\pi\epsilon_0 \frac{ab}{b-a}$$

This is inversely proportional to the resistance found in Exercise 1.4.

2.2

The planar capacitor formula is

$$C = \frac{\epsilon_0 A}{d};$$

Solving for A,

$$A = \frac{Cd}{\epsilon_0} = \frac{(1\text{ F})(10^{-9}\text{ m})}{10(8.85 \times 10^{-12}\text{ C}^2/\text{N}\cdot\text{m}^2)} = 1.1\text{ m}^2;$$

2.3

The solenoid inductor formula is

$$L = \frac{\mu_0 N^2 A}{l};$$

The loop area is $A = (5\text{ cm})^2 = 76\text{ cm}^2 = 7.6 \times 10^{-3}\text{ m}^2$. The density of coils is $N = (60)(0.05)(0.05\text{ m}) = 2.5 \times 10^4/\text{m}$.

Solving for l,

$$l = \frac{L}{\mu_0(Nl)^2} = \frac{1\text{ H}}{(7.6 \times 10^{-3}\text{ m}^2)(4 \times 10^7\text{ N/A}^2)(2.5 \times 10^4\text{ m})^2} = 17\text{ cm};$$

2.4

Integrate the flux over a rectangular section between the wires of length l. The magnetic field from a wire is

$$B = 2^{-\theta} \frac{I}{r} :$$

The flux from one wire is

$$= \int_0^a \frac{\mu_0 I}{2\pi r} dr = \frac{\mu_0 I}{2\pi} \ln(a=b)$$

The second wire contributes flux in the same direction, so the total flux is

$$\Phi_{\text{tot}} = \frac{\mu_0 I}{\pi} \ln(b=a)$$

The voltage drop is

$$V = \oint \mathbf{E} \cdot d\mathbf{l} = \frac{\mu_0 I}{4\pi} \ln(a=b)$$

Therefore

$$L = \frac{\mu_0 I}{V} \ln(a=b)$$

[Note: Technically we should account for the magnetic field inside each wire. The current inside the radius r is $I r^2/b^2$. $B(2r) = \mu_0 I r^2/b^2$, so

$$B = \frac{\mu_0 r I}{2b^2}$$

inside the wire. The integral of the flux inside the wire is

$$= \int_0^b \frac{\mu_0 r I}{2b^2} dr = \frac{\mu_0 I}{4}$$

The contribution to the inductance L from both wires is then

$$L^0 = \frac{\mu_0 I}{2}$$

which implies $L=L^0 = 2 \ln(a=b)$, which means the field inside the wire is negligible if $a \gg b$.]

The Biot-Savart law is

$$d\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \frac{I d\mathbf{l} \times \mathbf{R}}{R^3}$$

For a loop of radius L centered at $(0,0)$, the unit vector along $d\mathbf{l}$ is

$$\hat{\mathbf{t}} = -x\hat{\mathbf{i}} \sin\theta + y\hat{\mathbf{j}} \cos\theta$$

The distance from a point of the circle to a point r on the x-axis is

$$R = \sqrt{L^2 \cos^2\theta + L^2 \sin^2\theta + r^2}$$

The direction of \mathbf{R} is

$$\hat{\mathbf{R}} = \frac{L(\cos\theta \hat{\mathbf{i}} + \sin\theta \hat{\mathbf{j}}) - r\hat{\mathbf{i}}}{R} = \frac{1}{R} (\hat{\mathbf{x}}(L \cos\theta - r) + y\hat{\mathbf{y}} \sin\theta)$$

The cross product is

$$\mathbf{R} = \frac{z \hat{z}}{R(r \cos \theta - L)}$$

The integral of the Biot-Savart element is then

$$\mathbf{B} \sim \frac{\mu_0 I z \hat{z}}{4 \pi L^2} \frac{r \cos \theta - L}{L^2 [(L \cos \theta - r)^2 + L^2 \sin^2 \theta]^{3/2}}$$

This function involves elliptic integrals. Doing it numerically for $r = 0$ (the center of the loop), gives

$$\mathbf{B} \sim \frac{\mu_0 I z \hat{z}}{4 \pi L^2}$$

while at $r = L$, it is

$$\mathbf{B} \sim \frac{\mu_0 I z \hat{z}}{3 \pi L^2}$$

2.6

$$V_o = IR = RC \frac{dV_o}{dt} \quad (V - V_o) = RC \frac{dV_o}{dt}$$

This has the solution

$$V_o(t) = V e^{-t/RC}$$

2.7

Kirchho :

$$0 = RC \frac{dV_o}{dt} + V_o$$

This has the general solution

$$V_o(t) = V_1 e^{-t/RC} + V_2$$

Setting this to V at $t = t_0$ gives

$$V = V_1 e^{-t_0/RC} + V_2$$

The current through the capacitor is

$$I(t) = \frac{V}{R} = C \frac{dV_o}{dt} = \frac{V_1}{RC} e^{-t/RC}$$

which implies $V_1 = V e^{t_0/RC}$. Plugging this into the above gives $V_2 = 0$, so the solution is

$$V_o(t) = V e^{-(t-t_0)/RC}$$

This is decay to zero with the same time constant RC .

2.8

The integrator circuit of Figure 2.16(a) is governed by equation (2.4.13),

$$V_o + RC \dot{V}_o = V_i$$

For $V_o = \sin \omega t$, this implies

$$\sin \omega t + \omega RC \cos \omega t = V_i$$

When $\omega RC \gg 1$, the $\sin \omega t$ term is negligible, and V_i is proportional to \dot{V}_o , i.e., V_o is proportional to the antiderivative of V_i .

The differentiator circuit of Figure 2.16(b) is governed by equation (2.4.17),

For $V_o = \sin \omega t$, this implies

$$\omega \cos \omega t + \frac{1}{RC} \sin \omega t = V_i$$

When $\omega RC \gg 1$, the $\cos \omega t$ term is negligible, and V_i is proportional to \dot{V}_o , i.e., V_o is proportional to the derivative of V_i .

2.9

This is the circuit shown in Fig. 2.18(b). We have

$$\frac{V_o}{V_i} = \frac{R}{R + Z_C} = \frac{R}{R - j\omega C} = \frac{R}{R^2 + \omega^2 C^2} (R + j\omega C)$$

2.10

$$\frac{V_o}{V_i} = \frac{Z_L}{R + Z_L} = \frac{j\omega L}{R + j\omega L}$$

This is a high-pass filter.

2.11 From Exercise 2.9, we have

$$\frac{V_o}{V_i} = \frac{1}{1 - j\omega RC}$$

Solve for ω :

$$\omega = \frac{1}{RC} \sqrt{1 - |V_o/V_i|^2}$$

a)

$$3 \text{ dB} = 10 \log_{10} \frac{V_o}{V_i} = 10 \log_{10} \left(\frac{V_o}{V_i} \right)^2$$

$$\frac{V_o}{V_i} = 10^{0.3} = 1.96 \approx 2 \text{ at } t = 1 \text{ RC}$$

b)

$$\frac{V_o}{V_i} = 10^{0.1} = 1.26 \approx 1.3 \text{ at } t = 3 \text{ RC}$$

c)

$$\frac{V_o}{V_i} = 10^{0.2} = 1.58 \approx 1.6 \text{ at } t = 10 \text{ RC}$$

2.12

a)

$$10 \text{ dBm} = 20 \log_{10} \left(\frac{V}{V_0} \right)$$

$$\frac{V}{V_0} = 10^{1/2}$$

$$V = 3.16 V_0 = 1 \text{ V}$$

$$P = \frac{V^2}{2R} = 10 \text{ mW}$$

We could also have done this just by noting that 10 dBm is 10 times greater than 0 dBm. b)

$$35 \text{ dB} = 20 \log_{10} \frac{V_2}{V_1}$$

$$\frac{V_2}{V_1} = 10^{35/20} = 56$$

2.13

$$\sin \omega t = \frac{e^{i\omega t} - e^{-i\omega t}}{2i}$$

From (2.7.11),

$$F(\omega) = 2 \frac{e^{i\omega T} - e^{-i\omega T}}{2i}$$

2.14

f(t) is given by (note typo in book)

$$f(t) = \begin{cases} 0; & (n-1)T < t < nT \\ 1; & (n-1)T < t < (n+1)T \end{cases}$$

$$\begin{aligned}
 c_n &= \frac{1}{T} \int_{T=2}^T e^{i2nt-T} dt \\
 &= \frac{e^{i2n} - e^{-i2n}}{i2n} \\
 &= \frac{e^{3i n=2} - e^{-i n=2}}{i2n} = \frac{1}{n} e^{3i n=2} \sin(n=2)
 \end{aligned}$$

n	c _n
0	1/2
1	i
2	0
3	-i/3
4	0
5	i/5
6	0
7	-i/7

$$f(t) = \sum_{n=1}^{\infty} c_n e^{i2nt-T} = \frac{1}{2} + \sum_{n=1}^{\infty} \left(\frac{i}{n} e^{i2nt-T} - \frac{i}{n} e^{-i2nt-T} \right) = \frac{1}{2} + \sum_{n=1}^{\infty} \frac{2}{n} \sin(2nt-T)$$

The sum of the terms up through n = 7 (first seven nonzero terms) is shown in Figure 6.

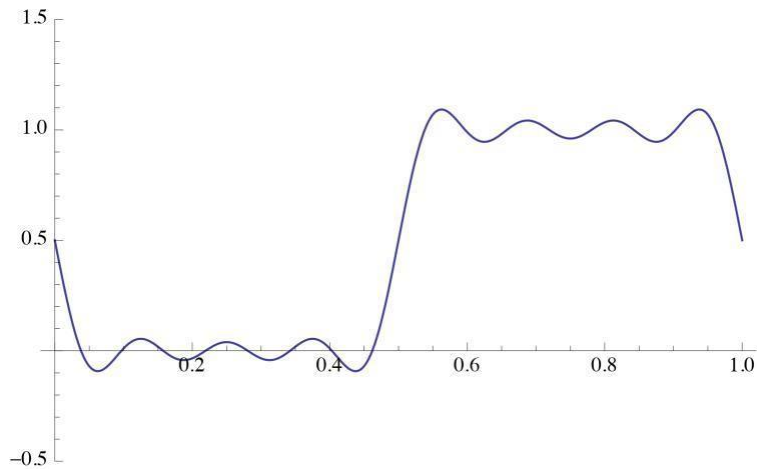


Figure 6: Fourier sum for Exercise 2.14.

2.15

The Fourier transform is

$$F(\omega) = \int_{-\infty}^{\infty} f(t) e^{i\omega t} dt$$

$f(t)$ is given by

$$f(t) = \begin{cases} 0; & (n-1)T < t < nT \\ 1; & nT < t < (n+1)T \end{cases}$$

The transform is

$$\begin{aligned} F(\omega) &= \sum_{n=-\infty}^{\infty} \int_{nT}^{(n+1)T} e^{i\omega t} dt \\ &= \sum_{n=-\infty}^{\infty} \left(\frac{e^{i\omega t}}{i\omega} \right) \Big|_{nT}^{(n+1)T} \\ &= \sum_{n=-\infty}^{\infty} \frac{e^{i\omega(n+1)T} - e^{i\omega nT}}{i\omega} \end{aligned}$$

The sum will be equal to infinity for $\omega = 2\pi n/T$ and zero otherwise (this is equivalent to a delta-function). Thus we have

$$F(\omega) = \begin{cases} \frac{1}{i\omega} & \omega = 2\pi n/T \\ 0 & \text{else} \end{cases}$$

This is the same as the result of Exercise 2.14.

2.16

The Fourier transform is

$$F(\omega) = \int_{-\infty}^{\infty} f(t) e^{i\omega t} dt = \int_0^{\infty} e^{-st} e^{i\omega t} dt = \frac{1}{s - i\omega}$$

The response function (2.5.7) is

$$\frac{V_0}{R} = \frac{i\omega C}{R} V_i = \frac{i\omega RC}{1 + i\omega RC} V_i$$

and the product is

$$F_o(\omega) = \frac{1}{1 + i\omega RC}$$

The reverse Fourier transform is

$$f(t) = \int_{-\infty}^{\infty} \frac{1}{1 + i\omega RC} e^{i\omega t} d\omega$$

For $t > 0$, $\int_{-\infty}^{\infty} \frac{1}{1 + i\omega RC} e^{i\omega t} d\omega$ converges, so this becomes

$$f(t) = \frac{1}{RC} \int_{-\infty}^{\infty} \frac{1}{i\omega + 1/RC} e^{i\omega t} d\omega = \frac{1}{RC} \int_{-\infty}^{\infty} \frac{1}{i\omega + 1/RC} e^{i\omega t} d\omega$$

Setting $s = 0$ gives

$$f(t) = 1 e^{-t/RC} :$$

2.17

Loops:

$$V_i = IR + I_C Z_C$$

$$V_i = IR + I_L Z_L$$

Node:

$$I = I_C + I_L$$

Solution:

$$I_C = \frac{Z_L}{R Z_C + R Z_L + Z_C Z_L} V_i = \frac{i \omega L}{i R \omega C + i \omega R L + L \omega C} V_i$$

$$V_C = I_C Z_C = \frac{L \omega C}{i R \omega C + i \omega R L + L \omega C} V_i$$

2.18

(2.8.2) is

$$\frac{V_o}{V_i} = \frac{R}{R + i \omega L + \frac{1}{i \omega C}} :$$

(2.8.12) is

$$V_i = IR + \frac{1}{C} \int I dt + LI$$

For $V_i = V_i(0)e^{i \omega t}$ and $I = I_0 e^{i \omega t}$, this becomes

$$i \omega V_i(0) e^{i \omega t} = i \omega I_0 R e^{i \omega t} + \frac{I_0}{C} e^{i \omega t} + L i \omega I_0 e^{i \omega t} :$$

The output V_o is across R , so $V_o = I_0 R e^{i \omega t}$. The above equation therefore becomes

$$i \omega V_i = i \omega V_o + \frac{V_o}{RC} + \frac{L \omega^2}{R} V_o$$

Solving for $V_o = V_i$ gives

$$\frac{V_o}{V_i} = \frac{i \omega R}{i \omega R + 1 + \omega^2 L R}$$

which is the same as (2.8.2).

2.19

We want a high-pass filter like that shown in Fig. 2.18(b), which has response (see Exercise 2.9)

$$\frac{V_o}{V_i} = \frac{1}{1 + j\omega RC}$$

We solve for C:

$$\frac{1}{1 + j\omega RC} = \frac{1}{\sqrt{2}} \quad \omega = 100 \text{ Hz} \quad \frac{1}{\sqrt{2}} = \frac{1}{\sqrt{1 + (\omega RC)^2}} \quad 1 = 0.6 F$$

2.20

Solve for 10% value:

$$e^{-t_{10}/\tau} = 0.1 \quad \tau = RC$$

For 90% value,

$$e^{-t_{90}/\tau} = 0.9$$

On the positive side,

$$t_{10} - t_{90} = \tau \ln(0.1) - \tau \ln(0.9) = \tau (\ln(0.1) - \ln(0.9)) = 1.2 \tau$$

2.21

The electric field is purely radial. For charge +Q on the inner sphere and -Q on the outer sphere, the electric field is

$$E(r) = \frac{Q}{4\pi\epsilon_0 r^2}$$

and the potential drop is

$$V(r) = \int_{r_1}^{r_2} E(r) dr = \frac{Q}{4\pi\epsilon_0 r_1} - \frac{Q}{4\pi\epsilon_0 r_2}$$

Comparing to the definition $Q = CV$, this gives

$$C = 4\pi\epsilon_0 \left(\frac{1}{r_1} - \frac{1}{r_2} \right)$$

2.22

The electric field from a line charge is found from Gauss's law,

$$E = \frac{\lambda}{2\pi\epsilon_0 r}$$

$$V = \int_a^b \frac{Q}{2\pi r} E(r) dr = \frac{Q}{2\pi} \ln(a=b):$$

By superposition, the other wire contributes the same, so we multiply by 2. This implies

$$C = \frac{2Q}{V} = \frac{2Q}{\frac{Q}{2\pi} \ln(a=b)} = \frac{4\pi}{\ln(a=b)}:$$

2.23

$$C = \frac{A}{d} = \frac{l w_0}{d}$$

$$L = \frac{\mu_0 l d}{w}$$

$$\frac{1}{LC} = \frac{1}{\mu_0 l^2 \frac{w_0}{d}} = \frac{d}{\mu_0 l^2 w_0} = \frac{c^2}{l^2} = 1:$$

2.24

The high-pass filter of Fig. 2.18(b) has response (see Exercise 2.9)

$$\frac{V_0}{V_i} = \frac{R}{R + i\omega C} = \frac{R}{R^2 + \omega^2 C^2}:$$

From (2.5.9),

$$\tan \phi = \frac{\text{Im } V_0/V_i}{\text{Re } V_0/V_i} = \frac{-\omega C}{R}$$

At high frequency, $\phi = 0$. At low frequency, $\phi = -90^\circ$.

2.25

$$\frac{V_0}{V_i} = \frac{j\omega L}{j\omega L + R} = \frac{j\omega L/R}{j\omega L/R + 1}$$

Figure ?? shows the plot with frequency in units of $R=L$.

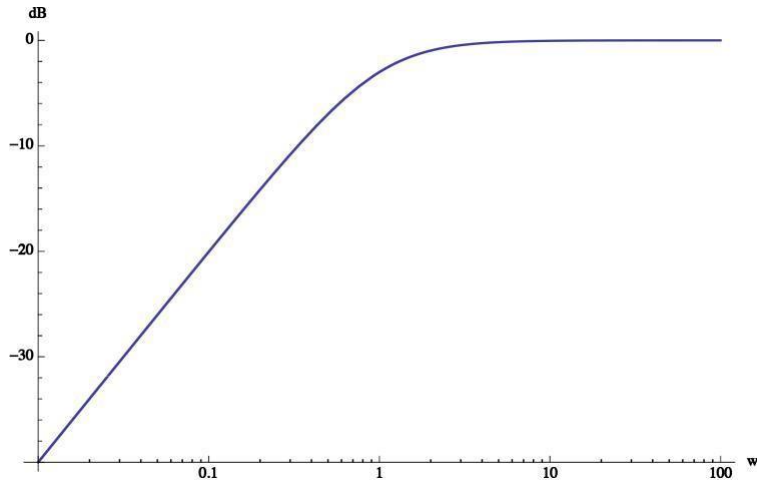


Figure 7: Response function for Exercise 2.25.

2.26

$|j\omega V_0|^2$ is increased by $(3/2)^2 = 2.25$.
 $10(\log_{10} 2.25) = 3.52$ dB

2.27

From Exercises 2.9 and 2.18:

$$\frac{V_0}{V_i} = \frac{1}{1 + j\omega RC}$$

Solving for ω ,

$$\omega = \frac{1}{RC} \sqrt{\frac{V_i^2}{V_0^2} - 1}$$

2.28

The low-pass filter response (2.5.7) is

$$\frac{V_0}{V_i} = \frac{1}{1 + j\omega RC}$$

Power response is

$$\frac{V_o}{V_i} = \frac{1}{\sqrt{1 + \omega^2 RC^2}} = \frac{1}{\sqrt{1 + \omega^2 RC^2}}$$

Solve for ω :

$$\omega = \frac{1}{RC}$$

10%:

$$\omega = \frac{3}{RC}$$

90%:

$$\omega = \frac{1}{3RC}$$

10-90 range:

$$\omega = \frac{2}{RC} \approx 667$$

2.29

$$\text{dB} = 20 \log_{10} \frac{V_2}{V_1}$$

$$\frac{V_2}{V_1} = 10^{\frac{\text{dB}}{20}} = 10^{\frac{13}{20}} = 4.5$$

The amplitude signal-to-noise ratio increases by a factor of 4.5, i.e. from 2 to 9.

It is also common to talk in terms of the signal-to-noise power ratio.

2.30

The circuit is shown in Figure 8.

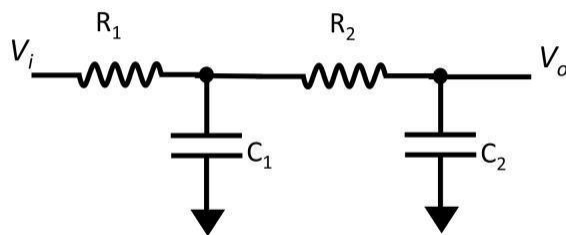


Figure 8: Circuit for Exercise 2.30.

Loops:

$$V_i = I_1 R_1 + I_{C1} Z_{C1}$$

$$V_i = I_1 R_1 + I_2 R_2 + V_o$$

$$V_o = I_2 Z_{C2}$$

Node:

$$I_1 = I_{C1} + I_2$$

Set $R_1 = R_2 = R$, $C_1 = C_2 = C$. Solution:

$$\frac{V_o}{V_i} = \frac{Z_C^2}{R^2 + 3RZ_C + Z_C^2}$$

In the Figure 9, the lower curve is this response function, while the upper curve is the single low-pass response, from Exercise 2.27.

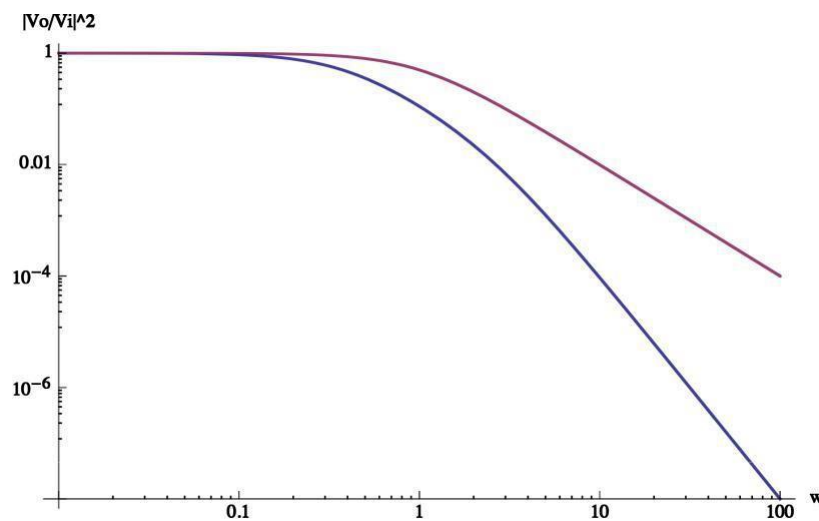


Figure 9: Response functions for Exercise 2.30.

2.31

Assuming perfect detector; no current flows into output; I_o is the current flowing from top to bottom in the right side of the circuit.

Loops:

$$V_i = I_1 Z_{C1} + I_2 R$$

$i \quad C1 \quad C$

R3 2

$$V_i = I_{R1}R + I_{C3} \frac{ZC}{2}$$

$$V_i = I_{R1}R + I_oR + V_o \frac{R}{2}$$

$$V_o = I_oZC + I_{R3} \frac{R}{2}$$

Nodes:

$$I_{R1} = I_{C3} + I_o$$

$$I_{C1} + I_o = I_{R3}$$

Solution:

$$\frac{V_o}{V_i} = \frac{(1 + C^2R^2 \omega^2)}{1 + 4iCR\omega + C^2R^2\omega^2}$$

This is plotted in Figure 10. When $\omega = 1/RC$, this equals 0. When $\omega \rightarrow 0$ or $\omega \rightarrow \infty$, it approaches unity.

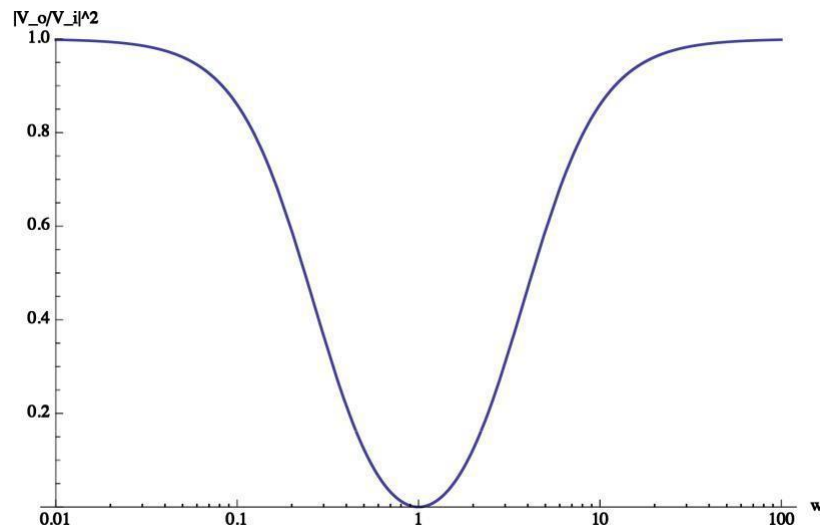


Figure 10: Response function for Exercise 2.31.

2.32

The circuit in the book is missing a 50- resistor between V_i and the rest of the circuit. As drawn, the circuit will give $V_o = V_i$ for all inputs. Putting a resistor there gives

Loops:

$$V_i = IR + V_o$$

$$V_o = I_1(ZC_1 + ZL_1)$$

$$V_o = I_2(Z_{C2} + Z_{L2})$$

Node:

$$I = I_1 + I_2$$

Solution:

$$\frac{V_o}{V_i} = \frac{(1 + C_1 L_1 \omega^2)(1 + C_2 L_2 \omega^2)}{1 + i C_1 R \omega + i C_2 R \omega + C_1 L_1 \omega^2 + C_2 L_2 \omega^2 + i C_1 C_2 L_1 R \omega^3 + i C_1 C_2 L_2 R \omega^3 + C_1 C_2 L_1 L_2 \omega^4}$$

This is a double notch filter, with zeroes where the two terms in the numerator vanish. Figure 11 shows a plot for the values given:

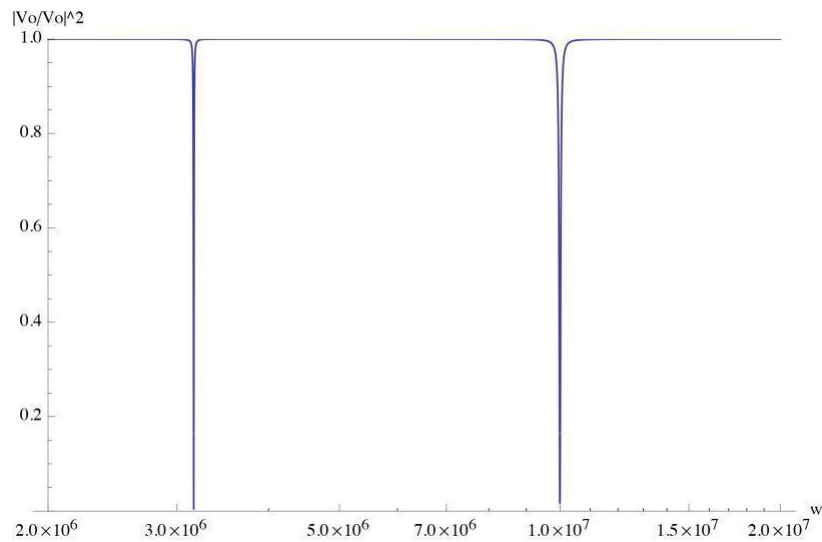


Figure 11: Double notch response function, for Exercise 2.32.

2.33

a) Capacitor relation:

$$I = C \frac{dV}{dt}$$

Loops:

$$V_s = I_1 R_1 + I_2 R_2$$

$$V_i = V_C + V_o \quad V_i = \frac{I}{C} + V_o$$

$$V_o = I_2 R_2$$

These become

$$V_s = I_1 R_1 + I_2 R_2$$

$$i!V_1 e^{i!t} = \frac{1}{C} \int I_C dt + i!V_{AC} e^{i!(t+)}$$

$$V_{DC} + V_{AC} e^{i!(t+)} = I_2 R_2$$

Node:

$$I_C + I_1 = I_2.$$

Write $I_1 = I_{1DC} + I_{1AC}$ and $I_2 = I_{2DC} + I_{2AC}$ (note that there is no DC component and AC terms equal: through the capacitor), and set DC terms equal

$$V_s = I_{1DC} R_1 + I_{2DC} R_2$$

$$0 = I_{1AC} R_1 + I_{2AC} R_2$$

$$i!V_1 e^{i!t} = \frac{1}{C} \int I_C dt + i!V_{AC} e^{i!(t+)}$$

$$V_{DC} = I_{2DC} R_2$$

$$V_{AC} e^{i!(t+)} = I_{2AC} R_2$$

$$I_C + I_{1AC} = I_{2AC}$$

$$I_{1DC} = I_{2DC}$$

We solve these for I_C ; I_{1DC} ; I_{1AC} ; I_{2DC} ; I_{2AC} ; V_{DC} ; and $V_{AC} e^i$, which gives

$$V_{DC} = \frac{R_2 V_s}{R_1 + R_2};$$

$$V_{AC} e^i = \frac{\frac{CR_1 R_2!}{iR_1 - iR_2 + CR_1 R_2!}}{i + CR_e!} V_1$$

$$= \frac{1}{i + CR_e!} V_1$$

where

$$R_e = \frac{R_1 R_2}{R_1 + R_2};$$

2.34

We want a low-pass filter that eliminates AC frequency of 60 Hz. We use the circuit shown in Fig. 2.18(a) with a polarized capacitor with the negative side grounded. Formula (2.5.11) gives us

$$jV_o = V_i j^2 = \frac{1}{1 + !^2 R^2 C^2};$$

We would like low series R to prevent DC droop of the voltage supply. Pick $R = 1 \Omega$, which is small compared to a typical 50 load impedance. Solve for C:

$$C = \frac{1}{R} \frac{q}{I_j} \frac{V_o - V_{ij}}{V_{ij}^2 - 1}$$

To eliminate 99% of the ripple, pick $\Delta V_o = V_{ij}^2 = 0.01$. Setting $\omega = 2\pi(60 \text{ Hz}) = 377 \text{ s}^{-1}$, we then have

$$C = \frac{1}{(377 \text{ s}^{-1})(50 \text{ V})} \Delta V_o = 0.00053 \text{ F} = 530 \text{ }\mu\text{F}$$